

2) Jackson 6.5

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a) $\vec{P}_{\text{field}} \equiv \int d^3x \vec{\nabla} \times \vec{B} = \epsilon_0 \mu_0 \int d^3x \vec{E} \times \vec{H}$



$= -\frac{1}{c^2} \int d^3x \vec{\nabla} \phi \times \vec{H}$, for static fields

$= -\frac{1}{c^2} \int d^3x [(\vec{\nabla} \times \phi \vec{H}) - \phi(\vec{\nabla} \times \vec{H})]$

$= -\frac{1}{c^2} \oint da \hat{r} \times (\phi \vec{H}) + \frac{1}{c^2} \int d^3x \phi(\vec{\nabla} \times \vec{H}) \quad (1)$

Provided that $\phi \vec{H} \sim \frac{1}{r^{2+\epsilon}}$ the surface integral vanishes

& for static fields $\vec{\nabla} \times \vec{H} = \vec{J}$

$\Rightarrow \vec{P}_{\text{field}} = \frac{1}{c^2} \int d^3x \phi \vec{J} \quad \#$

b) For, $\phi = \phi(t) - \vec{\nabla} \phi(t) \cdot \vec{x} = \phi_0 - \vec{E}_0 \cdot \vec{x}$

$\vec{P}_{\text{field}} = \frac{1}{c^2} \int d^3x (\phi_0 - \vec{E}_0 \cdot \vec{x}) \vec{J}(\vec{x})$

$= \frac{\phi_0}{c^2} \int d^3x \vec{J}(\vec{x}) - \frac{1}{c^2} \int \vec{E}_0 \cdot \vec{x} \vec{J}(\vec{x}) d^3x$

consider ith component:

$P_i = -\frac{1}{c^2} E_{0j} \int d^3x x_j J_i$

$= -\frac{E_{0j}}{c^2} \frac{1}{2} \int d^3x (x_j J_i - x_i J_j) = -\frac{E_{0j}}{c^2} \frac{1}{2} \int d^3x (-\epsilon_{ijk} \epsilon_{klm} x_l J_m)$

$= -\frac{E_{0j}}{c^2} (-\epsilon_{ijk}) \left(\frac{1}{2} \int d^3x \epsilon_{klm} x_l J_m \right) = \frac{E_{0j} \epsilon_{ijk}}{c^2} M_k$

$= \frac{1}{c^2} \epsilon_{ijk} E_{0j} M_k$

$\Rightarrow \vec{P} = \frac{1}{c^2} \vec{E}_0 \times \vec{m} \quad \# \quad (2)$

c) From multipole expansion:

$$\vec{A}(\vec{x}) = -\frac{\mu_0}{4\pi} \frac{\hat{r} \times \vec{m}}{r^2} + O\left(\frac{1}{r^3}\right)$$

$$\begin{aligned} \vec{H}(\vec{x}) &= \frac{1}{\mu_0} \vec{\nabla} \times \vec{A} \\ &= \frac{1}{4\pi} \left(\frac{3\hat{r}(\hat{r} \cdot \vec{m}) - \vec{m}}{r^3} \right) + O\left(\frac{1}{r^4}\right) \end{aligned} \quad (3)$$

Now consider the surface term from part (a) for

$$\phi = \phi_0 - \vec{E}_0 \cdot \vec{x}:$$

$$\begin{aligned} -\frac{1}{c^2} \oint da \hat{r} \times \phi \vec{H} &= -\frac{1}{c^2} \lim_{r \rightarrow \infty} \int d\Omega r^2 \hat{r} \times \vec{H}(\vec{x}) (\phi_0 - \vec{E}_0 \cdot \vec{x}) \\ &= -\frac{1}{c^2} \lim_{r \rightarrow \infty} \int d\Omega r^2 \hat{r} \times \left\{ \frac{1}{4\pi} \left(\frac{3\hat{r}(\hat{r} \cdot \vec{m}) - \vec{m}}{r^3} \right) + O\left(\frac{1}{r^4}\right) \right\} (\phi_0 - \vec{E}_0 \cdot \vec{x}) \\ &= -\frac{1}{c^2} \lim_{r \rightarrow \infty} \int d\Omega r^2 \hat{r} \times \left\{ \frac{1}{4\pi} \left(\frac{3\hat{r}(\hat{r} \cdot \vec{m}) - \vec{m}}{r^3} \right) + O\left(\frac{1}{r^4}\right) \right\} (\phi_0 - \vec{E}_0 \cdot \hat{r} r) \\ &= -\frac{1}{c^2} \lim_{r \rightarrow \infty} \int d\Omega \hat{r} \times \left\{ -\frac{(\vec{E}_0 \cdot \hat{r})}{4\pi} (3\hat{r}(\hat{r} \cdot \vec{m}) - \vec{m}) + O\left(\frac{1}{r}\right) \right\} \\ &= \frac{1}{4\pi c^2} \int d\Omega \hat{r} \times \{ 3\hat{r}(\hat{r} \cdot \vec{m}) - \vec{m} \} (\vec{E}_0 \cdot \hat{r}) \\ &= -\frac{1}{4\pi c^2} \int d\Omega \hat{r} \times \vec{m} \vec{E}_0 \cdot \hat{r} \\ &= -\frac{1}{4\pi c^2} \epsilon_{ijk} \hat{x}_i \int d\Omega r_j m_k E_{0l} r_l \\ &= -\frac{1}{4\pi c^2} \epsilon_{ijk} \hat{x}_i m_k E_{0l} \left(\frac{4\pi}{3} \delta_{jl} \right), \quad \int d\Omega r_i r_j = \frac{4\pi}{3} \delta_{ij} \\ &= -\frac{1}{3c^2} \epsilon_{ijk} \hat{x}_i m_k E_{0j} = -\frac{1}{3c^2} \vec{E}_0 \times \vec{m} \end{aligned}$$

Thus,

$$\vec{P}_{\text{field}} = \frac{1}{c^2} \dot{\vec{E}}_0 \times \vec{m} - \frac{1}{3c^2} \dot{\vec{E}}_0 \times \dot{\vec{m}}$$

$$= \frac{2}{3c^2} \dot{\vec{E}}_0 \times \dot{\vec{m}}, \quad \dot{\vec{E}}_0 \equiv \dot{\vec{E}}(0) \quad \#$$