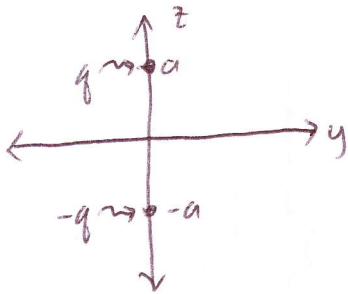


2) Jackson 3.6

ϕ/r



$$a) \phi(\vec{x}) = \frac{q}{4\pi\epsilon_0} \left(\frac{1}{|\vec{x} - a\hat{z}|} - \frac{1}{|\vec{x} + a\hat{z}|} \right) \quad (1)$$

note that:

$$\frac{1}{|\vec{x} - \vec{x}'|} = 4\pi \sum_{l=0}^{\infty} \sum_{m=-l}^l \frac{1}{2l+1} \frac{r_{<}^l}{r_{>}^{l+1}} Y_{lm}^*(\theta', \phi') Y_{lm}(\theta, \phi) \quad (2)$$

$$\rightarrow \frac{1}{|\vec{x} - a\hat{z}|} = 4\pi \sum_{l=0}^{\infty} \sum_{m=-l}^l \frac{1}{2l+1} \frac{r_{<}^l}{r_{>}^{l+1}} Y_{lm}^*(0, \phi') Y_{lm}(\theta, \phi) \quad (3)$$

$$\frac{1}{|\vec{x} + a\hat{z}|} = 4\pi \sum_{l=0}^{\infty} \sum_{m=-l}^l \frac{1}{2l+1} \frac{r_{<}^l}{r_{>}^{l+1}} Y_{lm}^*(\pi, \phi') Y_{lm}(\theta, \phi) \quad (4)$$

$$\text{but } Y_{lm}^*(\theta, \phi) = \sqrt{\frac{2l+1}{4\pi} \frac{(l-m)!}{(l+m)!}} P_l^m(\cos\theta) e^{-im\phi}$$

$$= \sqrt{\frac{2l+1}{4\pi} \frac{(l-m)!}{(l+m)!}} (-1)^m (1-x^2)^{m/2} \frac{d^m}{dx^m} P_l(1) e^{-im\phi}$$

$$= \begin{cases} 0 & \text{if } m \neq 0 \\ \sqrt{\frac{2l+1}{4\pi}} & , m = 0 \end{cases}, \text{ since } P_l(1) = \text{constant} = (1)^l \quad (5)$$

Similarly,

$$Y_{\ell m}^*(\pi, \varphi) = \sqrt{\frac{(2\ell+1)(\ell-m)!}{4\pi(\ell+m)!}} (-1)^m (1-x^2)^{m/2} \frac{d^m}{dx^m} P_{\ell}(-1) e^{-im\varphi}$$

$$= \begin{cases} 0 & \text{if } m \neq 0 \\ \sqrt{\frac{2\ell+1}{4\pi}} (-1)^{\ell} & \text{if } m=0 \end{cases} \quad (6)$$

(5), (6), (3), (4) \rightarrow (1):

$$\Phi(\vec{x}) = \frac{q}{4\pi\epsilon_0} \sum_{\ell=0}^{\infty} [1 - (-1)^{\ell}] \sqrt{\frac{4\pi}{2\ell+1}} \frac{r_{<}^{\ell}}{r_{>}^{\ell+1}} Y_{\ell 0}(\theta, \varphi)$$

$$= \frac{2q}{4\pi\epsilon_0} \sum_{\ell=0}^{\infty} \sqrt{\frac{4\pi}{2(2\ell+1)+1}} \frac{r_{<}^{2\ell+1}}{r_{>}^{(2\ell+1)+1}} Y_{(2\ell+1), 0}(\theta, \varphi)$$

$$= \frac{2q}{\sqrt{4\pi}\epsilon_0} \sum_{\ell=0}^{\infty} \sqrt{\frac{1}{4\ell+3}} \frac{r_{<}^{2\ell+1}}{r_{>}^{2\ell+2}} \left(\sqrt{\frac{2(2\ell+1)+1}{4\pi}} P_{2\ell+1}(\cos\theta) \right)$$

$$\Phi(\vec{x}) = \frac{q}{2\pi\epsilon_0} \sum_{\ell=0}^{\infty} \frac{r_{<}^{2\ell+1}}{r_{>}^{2\ell+2}} P_{2\ell+1}(\cos\theta) \quad (6)$$

b) $qa \equiv p/2 = \text{constant as } a \rightarrow 0$

since $a \rightarrow 0$, $r_{<} = a$ & $r_{>} = r$:

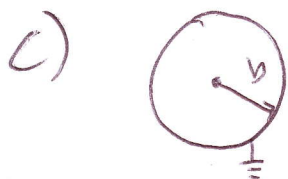
$$\Phi(\vec{x}) = \frac{q}{2\pi\epsilon_0} \sum_{\ell=0}^{\infty} \frac{a^{2\ell+1}}{r^{2\ell+2}} P_{2\ell+1}(\cos\theta)$$

$$= \frac{qa}{2\pi\epsilon_0 r^2} \sum_{\ell=0}^{\infty} \left(\frac{a}{r}\right)^{2\ell} P_{2\ell+1}(\cos\theta)$$

$$\phi(\vec{r}) = \frac{qa}{2\pi\epsilon_0 r^2} (P_1(\cos\theta) + O(a^2))$$

$$\approx \frac{qa}{2\pi\epsilon_0 r^2} P_1(\cos\theta), \text{ as } a \rightarrow 0$$

$$\phi(\vec{r}) = \frac{P \cos\theta}{4\pi\epsilon_0 r^2} \quad \checkmark \quad (7)$$



at $r < b$:

$$\phi_{\text{shell}} = \sum_{n=0}^{\infty} \left(\frac{a_n}{r^{n+1}} + b_n r^n \right) P_n(\cos\theta) \quad \checkmark$$

$$\phi_{\text{shell}}(0) = \text{Finite}$$

$$\rightarrow a_n = 0 \quad \checkmark$$

$$\phi = \phi_{\text{dipole}} + \phi_{\text{shell}}$$

$$= \frac{P \cos\theta}{4\pi\epsilon_0 r^2} + \sum_{n=0}^{\infty} b_n r^n P_n(\cos\theta) \quad \checkmark \quad (8)$$

at $r = b$:

$$\phi(b) = 0$$

$$\frac{P \cos\theta}{4\pi\epsilon_0 b^2} + b_0 + b_1 b \cos\theta + \sum_{n=2}^{\infty} b_n r^n P_n(\cos\theta) = 0$$

$$b_0 + \left(\frac{P}{4\pi\epsilon_0 b^2} + b_1 b \right) \cos\theta + \sum_{n=2}^{\infty} b_n r^n P_n(\cos\theta) = 0$$

$P_n(\cos\theta)$ are linearly independent functions thus;

$$b_n = 0, n \neq 1$$

$$b_1 b + \frac{P}{4\pi\epsilon_0 b^2} = 0$$

$$b_1 = \frac{-P}{4\pi\epsilon_0 b^3} \quad (a)$$

(a) \rightarrow (8):

$$\phi = \frac{p \cos\theta}{4\pi\epsilon_0 r^2} - \frac{P}{4\pi\epsilon_0 b^3} r \cos\theta$$

$$= \frac{P}{4\pi\epsilon_0} \left(\frac{1}{r^2} - \frac{r}{b^3} \right) \cos\theta \quad \checkmark \quad \#$$